# Tsirelson's problem and Kirchberg's conjecture 

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## Overview

- Quantum-mechanical axioms of composite systems.
- Their significance for nonlocality theory. (Tsirelson's problem)
- Strong ties to an open problem on $C^{*}$-algebras. (Kirchberg's conjecture)
- Two new paradigms of quantum correlations.

Disclaimer: Some of these results were obtained independently in
Junge, Navascués, Palazuelos, Pérez-García, Scholz, Werner,
Connes' embedding problem and Tsirelson's problem, J. Math. Phys. 52, 012102 (2011).

## Composite systems in quantum theory I

## Tensor product assumption:

- The state space of a joint system composed out of two subsystems is a tensor product

$$
\mathcal{H}_{A} \otimes \mathcal{H}_{B}
$$

with local observables

$$
A \otimes \mathbb{1} \text { for } A \in \mathcal{B}\left(\mathcal{H}_{A}\right), \quad \mathbb{1} \otimes B \text { for } B \in \mathcal{B}\left(\mathcal{H}_{B}\right) .
$$

- Posited by standard quantum theory.
- The composite system can be constructed from the subsystems.


## Composite systems in quantum theory II

## Commutativity assumption:

- State space of a joint system is a Hilbert space $\mathcal{H}$ with local observables

$$
A, B \in \mathcal{B}(\mathcal{H}) \text { such that } A B=B A
$$

- The joint state space is in general not uniquely determined by the subsystems.
- Defining $\mathcal{H} \equiv \mathcal{H}_{A} \otimes \mathcal{H}_{B}$ obviously works $\Rightarrow$ tensor product assumption is a special case.
- For finite-dimensional systems, it is essentially equivalent to the tensor product assumption.
- Not so in infinite dimensions!
( $\rightarrow$ This is what the theory of $C^{*}$-tensor products is about.)


## Composite systems in quantum theory III

Philosophical observation:

- Nature does not construct composite systems from subsystems.
- Rather, she presents us composite systems which we perceive as made out of subsystems.
- $\Rightarrow$ Correct question: when does a system look like it were composed out of subsystems?
- One possible answer: the commutativity assumption.


## Nonlocality theory: Bell scenarios



- with the tensor product assumption:

$$
\begin{gathered}
A_{a}^{x} \in \mathcal{B}\left(\mathcal{H}_{A}\right), \quad B_{b}^{y} \in \mathcal{B}\left(\mathcal{H}_{B}\right), \quad \psi \in \mathcal{H}_{A} \otimes \mathcal{H}_{B} \\
P(a, b \mid x, y)=\left\langle\psi,\left(A_{a}^{x} \otimes B_{b}^{y}\right) \psi\right\rangle
\end{gathered}
$$

## Nonlocality theory: Bell scenarios



- with the commutativity assumption:

$$
\begin{gathered}
A_{a}^{x}, B_{b}^{y} \in \mathcal{B}(\mathcal{H}), \quad A_{a}^{x} B_{b}^{y}=B_{b}^{y} A_{a}^{x}, \quad \psi \in \mathcal{H} \\
P(a, b \mid x, y)=\left\langle\psi, A_{a}^{x} B_{b}^{y} \psi\right\rangle
\end{gathered}
$$

## Sets of quantum correlations

- Two sets of quantum correlations $P(a, b \mid x, y) \in \mathbb{R}^{k^{2} m^{2}}$,

1. With the tensor product assumption: $\mathcal{Q}_{\otimes}$.
2. With the commutativity assumption: $\mathcal{Q}_{c}$.

- $\mathcal{Q}_{\otimes}$ is convex, $\mathcal{Q}_{c}$ is closed convex.
- It is unclear whether $\mathcal{Q}_{\otimes}$ is closed; consider the closure $\overline{\mathcal{Q}}_{\otimes}$.
- Tensor product assumption is a special case of the commutativity assumption $\Rightarrow \overline{\mathcal{Q}}_{\otimes} \subseteq \mathcal{Q}_{c}$.
- Tsirelson's problem: $\overline{\mathcal{Q}}_{\otimes} \stackrel{?}{=} \mathcal{Q}_{c}$


## More on Tsirelson's problem

Tsirelson's problem: $\overline{\mathcal{Q}}_{\otimes} \stackrel{?}{=} \mathcal{Q}_{c}$

- The answer may depend on the Bell scenario considered.
- Every $P(a, b \mid x, y) \in \mathcal{Q}_{c}$ coming from a state with finite-dimensional $\mathcal{H}$ lies also in $\mathcal{Q}_{\otimes}$.
- No further results are known.
- Physical relevance of a potential negative answer is unclear.

Additional motivation from nonlocality theory:

- Most (all?) examples of quantum correlations use the tensor product assumption.
- Most (all?) upper bounds on quantum correlations use the commutativity assumption. (E.g. the semidefinite hierarchy.)
$\rightarrow$ Tsirelson's problem: will improvement of lower and upper bounds lead to convergence, or will there remain a gap?


## Kirchberg's QWEP conjecture

Let $\mathbb{F}_{2}$ be the free group on two generators. For a discrete group $G$, $C^{*}(G)$ denotes its maximal group $C^{*}$-algebra.

- Kirchberg's QWEP conjecture:

$$
C^{*}\left(\mathbb{F}_{2}\right) \otimes_{\max } C^{*}\left(\mathbb{F}_{2}\right) \stackrel{?}{=} C^{*}\left(\mathbb{F}_{2}\right) \otimes_{\min } C^{*}\left(\mathbb{F}_{2}\right)
$$

- Proposed by Kirchberg in $1993^{1}$.
- Equivalent to Connes' embedding problem from 1976.
- Many reformulations exist as questions on $C^{*}$-algebras, von Neumann algebras, operator spaces...
${ }^{1}$ E. Kirchberg, On nonsemisplit extensions, tensor products and exactness of group C*-algebras, Invent. Math. (1993).


## From Kirchberg's conjecture to Tsirelson's problem

Theorem
If QWEP is true, then $\overline{\mathcal{Q}}_{\otimes}=\mathcal{Q}_{c}$ in all Bell scenarios.

The proof relies on a characterization of both $\overline{\mathcal{Q}}_{\otimes}$ and $\mathcal{Q}_{c}$ in terms of group $C^{*}$-algebras.

## Quantum correlations and group $C^{*}$-algebras

Sketch of the connection to group $C^{*}$-algebras:

- Label the outcomes of an m-outcome measurement by the $m$ th roots of unity $e^{2 \pi i \frac{j}{m}}$.
$\Rightarrow$ Observable is a unitary operator of order $m$.
- $k m$-outcome measurements correspond to $k$ unitaries of order $m$.
- This is equivalent to a unitary representation of the discrete group

$$
\Gamma=\underbrace{\mathbb{Z}_{m} * \ldots \mathbb{Z}_{m}}_{k \text { factors }}
$$

and hence to a representation of $C^{*}(\Gamma)$.

- This allows the use of methods from $C^{*}$-algebra theory, group theory, and representation theory.
- Example application: CHSH is simple because $\mathbb{Z}_{2} * \mathbb{Z}_{2} \cong \mathbb{Z} \rtimes \mathbb{Z}_{2}$.


## Spatiotemporal correlations I

- In a usual Bell scenario, each party conducts exactly one measurement per run.
- More generally, one can consider scenarios where each party is allowed to conduct several measurements per run in temporal succession. $\rightarrow$ "Spatiotemporal correlations"
- Motivated by the group $C^{*}$-algebra approach to quantum correlations.
- There are examples of $A_{a}^{x}$ and $B_{B}^{y}$ and an initial state $\psi$ such that the ensuing spatial correlations are local, but the spatiotemporal correlations prove nonlocality.


## Spatiotemporal correlations II

Tsirelson's problem can be formulated analogously for spatiotemporal correlations.

Theorem
If QWEP is true, then Tsirelson's problem for spatiotemporal correlations also has a positive answer.

The proof is exactly analogous to the purely spatial case.

## Steering data I

- If Alice conducts measurement $x$ and gets the outcome $a$, then Bob's system collapses to the state

$$
\rho(a \mid x)=\operatorname{tr}_{A}\left(\rho\left(A_{a}^{x} \otimes \mathbb{1}\right)\right) \quad(\text { unnormalized: } \operatorname{tr}(\rho(a \mid x))=P(a \mid x))
$$

- "Steering" of Bob's system by Alice.
- We call the set of unnormalized states $\rho(a \mid x)$ steering data.
- Steering data is a quantum analogue of the conditional probability distribution $P(a \mid x)$ :

$$
\text { classical: } P(a \mid x) \rightsquigarrow \text { quantum: } \rho(a \mid x)
$$

## Theorem

Steering data $\rho(a \mid x)$ can arise in this way if and only if it satisfies the no-signaling condition:

$$
\sum_{a} \rho(a \mid x) \text { is independent of } x
$$

## Bipartite steering data I

- Similarly, one can consider the case where both Alice and Bob measure and can steer the system of a third party:
$\rho(a \mid x)=\operatorname{tr}_{A B}\left(\rho\left(A_{a}^{x} \otimes \mathbb{1} \otimes \mathbb{1}\right)\right), \quad \rho(b \mid y)=\operatorname{tr}_{A B}\left(\rho\left(\mathbb{1} \otimes B_{b}^{y} \otimes \mathbb{1}\right)\right)$.
- The system of the third party is taken to be of fixed dimension $d$, so that $\rho(a \mid x), \rho(b, y) \in M_{d}(\mathbb{C})$.
- This is bipartite steering data.
- No joint measurements needed: $\rho(a \mid x)$ and $\rho(b \mid y)$ are quantum analogues of the marginals $P(a \mid x)$ and $P(b \mid y)$ !
- Again, there is an obvious variant of Tsirelson's problem.


## Bipartite steering data II

## Theorem

1. If QWEP holds, then Tsirelson's problem for bipartite steering data has a positive answer in all Bell scenarios.
2. If Tsirelson's problem for bipartite steering data has a positive answer in some (non-CHSH) Bell scenario, then QWEP holds.
$\rightarrow$ Tsirelson's problem for bipartite steering data can be considered a physical reformulation of the QWEP conjecture.

## Diagram of implications

Bipartite Bell scenario with $k$ measurements per party and $m$ outcomes each.


## Summary

- We replace the tensor product axiom of quantum mechanics by the commutativity assumption and study whether this allows more nonlocality ("Tsirelson's problem").
- Assuming the validity of Kirchberg's QWEP conjecture, we find a positive answer to this. Likewise for Tsirelson's problem on spatiotemporal quantum correlations.
- The proof is based on relating quantum correlations to group $C^{*}$-algebras.
- Tsirelson's problem for bipartite steering data is equivalent to QWEP for every (non-CHSH) Bell scenario.
- In particular, a proof or counterexample in any scenario would at the same time decide the problem for all other scenarios.


## Backup slides

## Proving nonlocality by spatiotemporal correlations I

- 2 qubits on Alice's side, 1 qubit on Bob's side. W state:

$$
|W\rangle=\frac{1}{\sqrt{3}}|00\rangle \otimes|1\rangle+\frac{1}{\sqrt{3}}(|01\rangle+|10\rangle) \otimes|0\rangle
$$

- $\pm 1$-valued observables:

$$
\begin{gathered}
A_{1}=\sigma_{z} \otimes \mathbb{1}, \quad A_{2}=\sigma_{x} \otimes \mathbb{1}, \quad A_{3}=\mathbb{1} \otimes \sigma_{z}, \quad A_{4}=\mathbb{1} \otimes \sigma_{x}, \\
B_{1}=\frac{\sigma_{z}-\sigma_{x}}{\sqrt{2}}, \quad B_{2}=\frac{\sigma_{z}+\sigma_{x}}{\sqrt{2}} .
\end{gathered}
$$

## Proving nonlocality by spatiotemporal correlations II

- If Alice measures only once: $A_{3}$ and $A_{4}$ are redundant, since they yield the same joint statistics as $A_{1}$ and $A_{2}$.
$\Rightarrow$ Situation is equivalent to a 2 -qubit system in the CHSH scenario.
By calculation: no CHSH violation.
- When Alice is allowed to measure twice: she begins with $A_{1}$ and then chooses between $A_{3}$ and $A_{4}$. The ensuing correlations show a Hardy-type nonlocality.

It's a bad example: Alice's sequential measurements commute.
Open problems:

- Find better examples!
- Are there examples where the correlations are local with up to $n$ sequential measurements, but nonlocal with $n+1$ sequential measurements?


## The multipartite Tsirelson's problem

- Tsirelson's problem, its two extensions, and the QWEP conjecture generalize all to the multipartite case.
- It seems that the same considerations as in the bipartite case imply the same results; details will have to checked.
- Even if the usual QWEP conjecture is true, the multipartite Tsirelson's problem may still be false.

